

## 2024年《电磁学》期中考试试题解答

1. (25分) 一个半径为  $R$  的均匀介质球, 相对介电常数为  $\epsilon_r$ , 球内均匀分布有总电量为  $Q_0$  的自由电荷, 球外为真空。求: (1) 球体内的极化电荷和球面的极化电荷; (2) 球内的静电能、球外的静电能和总静电能。(3) 该系统的宏观静电能为多少? 极化能为多少? (4) 如果该介质球的电导率为  $\sigma$ , 求任意时刻  $t$  球内  $r$  处的电场强度、电流密度; (5) 整个漏电过程的球内产生的总焦耳热为多少?

【解】(1) 由高斯定理, 球内的电场强度为:

$$\vec{E} = \frac{\rho_0}{3\epsilon_0\epsilon_r} \vec{r}$$

极化强度为:

$$\vec{P} = \epsilon_0(\epsilon_r - 1)\vec{E} = \frac{(\epsilon_r - 1)\rho_0}{3\epsilon_r} \vec{r}$$

极化电荷体密度为:

$$\rho' = -\nabla \cdot \vec{P} = -\frac{(\epsilon_r - 1)\rho_0}{3\epsilon_r} \nabla \cdot \vec{r}$$

因为  $\nabla \cdot \vec{r} = \frac{\partial x}{\partial x} + \frac{\partial y}{\partial y} + \frac{\partial z}{\partial z} = 3$ , 所以:

$$\rho' = -\frac{(\epsilon_r - 1)\rho_0}{\epsilon_r}$$

球内总极化电荷为:

$$Q'_v = \frac{4}{3}\pi R^3 \rho' = -\frac{4\pi R^3 (\epsilon_r - 1)\rho_0}{3\epsilon_r} = -\frac{(\epsilon_r - 1)}{\epsilon_r} Q_0$$

球表面极化电荷面密度:

$$\sigma' = P_n = \frac{(\epsilon_r - 1)\rho_0}{3\epsilon_r} R$$

球表面极化电荷

$$Q'_s = 4\pi R^2 \sigma' = \frac{4\pi R^3 (\epsilon_r - 1)\rho_0}{3\epsilon_r} = \frac{(\epsilon_r - 1)}{\epsilon_r} Q_0$$

球的总极化电荷

$$Q' = Q'_v + Q'_s = 0$$

(2) 球内的静电能为:

$$W_{\text{静球内}} = \iiint_V \frac{1}{2} \varepsilon_0 \varepsilon_r E^2 dV = \frac{1}{2} \varepsilon_0 \varepsilon_r 4\pi \int_0^R \left( \frac{\rho_0}{3\varepsilon_0 \varepsilon_r} r \right)^2 r^2 dr = \frac{2\pi \rho_0^2}{9\varepsilon_0 \varepsilon_r} \frac{1}{5} R^5 = \frac{Q_0^2}{40\pi \varepsilon_0 \varepsilon_r R}$$

球外任一点的电场强度为:

$$\vec{E} = \frac{Q_0}{4\pi \varepsilon_0 r^3} \vec{r}$$

球外的静电能为:

$$W_{\text{静球外}} = \iiint_V \frac{1}{2} \varepsilon_0 E^2 dV = \frac{1}{2} \varepsilon_0 4\pi \int_R^\infty \left( \frac{Q_0}{4\pi \varepsilon_0 r^2} \right)^2 r^2 dr = \frac{Q_0^2}{8\pi \varepsilon_0} \int_R^\infty \frac{1}{r^2} dr = \frac{Q_0^2}{8\pi \varepsilon_0 R}$$

总静电能为:

$$W_{\text{静}} = W_{\text{静球内}} + W_{\text{静球外}} = \frac{Q_0^2}{40\pi \varepsilon_0 \varepsilon_r R} + \frac{Q_0^2}{8\pi \varepsilon_0 R} = \frac{Q_0^2}{8\pi \varepsilon_0 R} \left( 1 + \frac{1}{5\varepsilon_r} \right)$$

(3) 宏观静电能密度为:

$$w_{e0} = \frac{1}{2} \varepsilon_0 E^2$$

因为球外  $D = \varepsilon_0 E$ , 因此球外的宏观静电能就是静电能:

$$W_{e0\text{球外}} = \frac{Q_0^2}{8\pi \varepsilon_0} \int_R^\infty \frac{1}{r^2} dr = \frac{Q_0^2}{8\pi \varepsilon_0 R}$$

球内的宏观静电能为:

$$W_{e0\text{球内}} = \iiint_V \frac{1}{2} \varepsilon_0 E^2 dV = \frac{1}{2} \varepsilon_0 4\pi \int_0^R \left( \frac{\rho_0}{3\varepsilon_0 \varepsilon_r} r \right)^2 r^2 dr = \frac{2\pi \rho_0^2}{9\varepsilon_0 \varepsilon_r} \frac{1}{5} R^5 = \frac{Q_0^2}{40\pi \varepsilon_0 \varepsilon_r^2 R}$$

总宏观静电能为:

$$W_{e0} = W_{e0\text{球内}} + W_{e0\text{球外}} = \frac{Q_0^2}{40\pi \varepsilon_0 \varepsilon_r^2 R} + \frac{Q_0^2}{8\pi \varepsilon_0 R} = \frac{Q_0^2}{8\pi \varepsilon_0 R} \left( 1 + \frac{1}{5\varepsilon_r^2} \right)$$

因为:

$$W_{\text{静}} = W_{e0} + W_{\text{极}}$$

所以极化能为:

$$W_{\text{极}} = W_{\text{静}} - W_{e0} = \frac{Q_0^2}{8\pi \varepsilon_0 R} \left( 1 + \frac{1}{5\varepsilon_r} \right) - \frac{Q_0^2}{8\pi \varepsilon_0 R} \left( 1 + \frac{1}{5\varepsilon_r^2} \right) = \frac{Q_0^2}{40\pi \varepsilon_0 \varepsilon_r R} \left( 1 - \frac{1}{\varepsilon_r} \right) = \frac{Q_0^2 (\varepsilon_r - 1)}{40\pi \varepsilon_0 \varepsilon_r^2 R}$$

验算(非必须):

$$\begin{aligned} W_{\text{极}} &= \frac{1}{2} \iiint_V \vec{P} \cdot \vec{E} dV = \frac{1}{2} \frac{(\varepsilon_r - 1) \rho_0}{3\varepsilon_r} \cdot \frac{\rho_0}{3\varepsilon_0 \varepsilon_r} \int_0^R r^2 4\pi r^2 dr = \frac{2\pi (\varepsilon_r - 1) \rho_0^2}{9\varepsilon_0 \varepsilon_r^2} \frac{1}{5} R^5 \\ &= \frac{2\pi (\varepsilon_r - 1) 9Q_0^2}{9\varepsilon_0 \varepsilon_r^2 (4\pi R^3)^2} \frac{1}{5} R^5 = \frac{(\varepsilon_r - 1) Q_0^2}{40\pi \varepsilon_0 \varepsilon_r^2 R} \end{aligned}$$

(4) 设任意时刻球内的自由电荷密度的为  $\rho(t)$ , 则由斯定理得到电场强度为:

$$\vec{E}(t) = \frac{\rho(t)}{3\epsilon_0\epsilon_r} \vec{r}$$

半径为  $r$  的球内的自由电荷总量为:

$$Q_r(t) = \frac{4}{3}\pi r^3 \rho(t)$$

由电荷守恒方程, 有:

$$4\pi r^2 j(t) = -\frac{dQ_r}{dt} = -\frac{4}{3}\pi r^3 \frac{d\rho}{dt}$$

$$4\pi r^2 \sigma E = -\frac{4\pi r^3}{3} \frac{d\rho}{dt}$$

$$\frac{d\rho}{\rho} = -\frac{\sigma}{\epsilon_0\epsilon_r} dt$$

积分得到 (利用初始条件,  $t=0$  时,  $\rho=\rho_0$ ):

$$\rho = \rho_0 e^{-\frac{\sigma}{\epsilon_0\epsilon_r} t} = \frac{3Q_0}{4\pi R^3} e^{-\frac{\sigma}{\epsilon_0\epsilon_r} t}$$

因此球内任意时刻的电场强度为:

$$\vec{E}(t) = \frac{\rho(t)}{3\epsilon_0\epsilon_r} \vec{r} = \frac{Q_0}{4\pi\epsilon_0\epsilon_r R^3} e^{-\frac{\sigma}{\epsilon_0\epsilon_r} t} \vec{r}$$

电流体密度为:

$$\vec{j}(t) = \sigma \vec{E}(t) = \frac{\sigma Q_0}{4\pi\epsilon_0\epsilon_r R^3} e^{-\frac{\sigma}{\epsilon_0\epsilon_r} t} \vec{r}$$

(5) 球内自由电荷最后全部流到外球面上, 焦耳功率密度为:

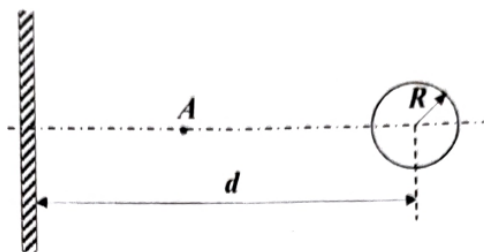
$$p = \sigma E^2 = \frac{\sigma Q_0^2}{(4\pi\epsilon_0\epsilon_r R^3)^2} e^{-\frac{2\sigma}{\epsilon_0\epsilon_r} t} r^2$$

总焦耳热为:

$$\begin{aligned} W &= \int_0^\infty dt \iiint_V p dV = \frac{\sigma Q_0^2}{(4\pi\epsilon_0\epsilon_r R^3)^2} \int_0^\infty e^{-\frac{2\sigma}{\epsilon_0\epsilon_r} t} dt \int_0^R r^2 4\pi r^2 dr \\ &= \frac{\sigma Q_0^2}{(4\pi\epsilon_0\epsilon_r R^3)^2} \left( -\frac{\epsilon_0\epsilon_r}{2\sigma} \right) e^{-\frac{2\sigma}{\epsilon_0\epsilon_r} t} \Big|_0^\infty \left[ \frac{4\pi}{5} r^5 \Big|_0^R \right] \\ &= \frac{\sigma Q_0^2}{(4\pi\epsilon_0\epsilon_r R^3)^2} \frac{\epsilon_0\epsilon_r}{2\sigma} \frac{4\pi}{5} R^5 = \frac{Q_0^2}{40\pi\epsilon_0\epsilon_r R} \end{aligned}$$

焦耳热正好等于初始时刻球内的静电能。

2. (25分) 在扫描隧道显微镜 (STM) 实验测量中, 可以近似地把 STM 探头看成一半径为  $a$  ( $nm$  量级) 的导体球, 待测样品 ( $mm$  量级) 可以看成是一块无限大接地导体平面, 设球心与样品相距为  $d$ , 并设  $d \gg a$ .



求:

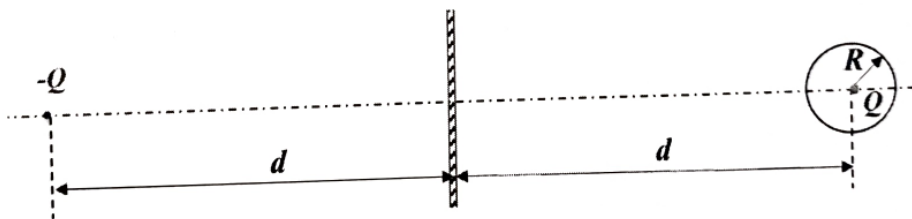
- (1) 探头和样品之间电容的零阶近似值;
- (2) 探头和样品之间电容的一阶近似值(点电荷模型); 这时探头受到的作用力为多大?
- (3) 探头和样品之间电容的二阶近似值 (电偶极子模型, 已知导体球在外场  $E_0$  中感应电偶极矩为  $p=4\pi\epsilon_0 a^3 E_0$ );
- (4) 在 (3) 近似下, 探头受到作用力为多大? (电偶极子受力用  $\vec{F} = (\vec{p} \cdot \nabla) \vec{E}$  计算)
- (5) 当探头带电量为  $Q$  时, 在 (3) 的近似下, 将探头与样品完全分离需提供多少的能量?

5' 【解】(1) 零阶近似, 即可以认为球与导体平面相距无限远, 即为孤立导体球的电容:

$$C_0 = 4\pi\epsilon_0 a$$

6' (2) 一阶近似, 设导体球带电量为  $Q$ , 在导体平面的感应电荷当成一个像, 电量  $q' = -Q$ , 位置  $d' = -d$ , 导体球面的电势为:

$$U = \frac{Q}{4\pi\epsilon_0 a} + \frac{-Q}{4\pi\epsilon_0 2d} = \frac{Q}{4\pi\epsilon_0} \left( \frac{1}{a} - \frac{1}{2d} \right)$$



电容值为:

$$C_1 = \frac{Q}{U} = \frac{4\pi\epsilon_0}{\left( \frac{1}{a} - \frac{1}{2d} \right)} = \frac{8\pi\epsilon_0 ad}{(2d - a)}$$

因为  $d \gg a$ , 也可以简化为 (这一步非必须):

$$C_1 = \frac{8\pi\epsilon_0 ad}{2d(1 - a/2d)} = 4\pi\epsilon_0 a(1 + a/2d) = 4\pi\epsilon_0 a + 2\pi\epsilon_0 a^2/d = C_0 + \Delta C_1$$

$$\Delta C_1 = 2\pi\epsilon_0 a^2/d$$

探头的作用力为：

$$F = -\frac{1}{4\pi\epsilon_0} \frac{Q^2}{4d^2} = -\frac{1}{16\pi\epsilon_0} \frac{Q^2}{d^2}$$

负号表示吸引力。

(3) 二阶近似，采用电偶极子模型：即像电荷 $-Q$ 在导体球周围产生的电场视为均匀电场：

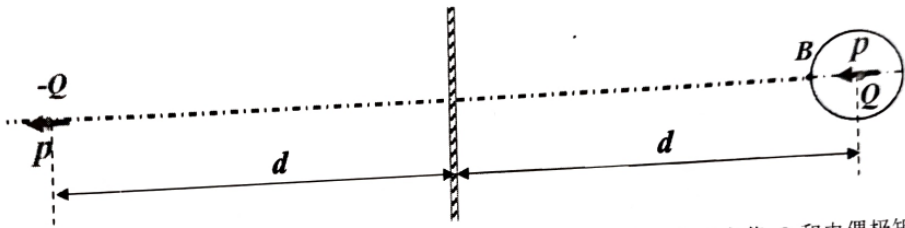
$$E_0 = \frac{Q}{16\pi\epsilon_0 d^2}$$

导体球面的感应电荷等效于一个放置在球心的电偶极矩，大小为：

$$p = 4\pi\epsilon_0 a^3 E_0 = \frac{Qa^3}{4d^2}$$

该电偶极子进一步在无限大导体左侧成一个像，像电偶极矩大小为：

$$p' = p$$



最终，导体球等效一个球心处的点电荷 $Q$ 和电偶极矩 $p$ ，像电荷也是点电荷 $-Q$ 和电偶极矩 $p'$ ，如图所示。球面的总电势为：

$$U = \frac{Q}{4\pi\epsilon_0 a} + \frac{-Q}{4\pi\epsilon_0 2d} + \frac{p \cos \theta_1}{4\pi\epsilon_0 R^2} + \frac{p' \cos \theta_2}{4\pi\epsilon_0 (2d)^2}$$

作为近似，取球面 $B$ 点计算电势（等同于取球心计算，因为 $d \gg a$ ），即有：

$$U = \frac{Q}{4\pi\epsilon_0 a} + \frac{-Q}{4\pi\epsilon_0 2d} + \frac{p}{4\pi\epsilon_0 a^2} - \frac{p}{4\pi\epsilon_0 (2d)^2} \quad \text{球心, } U = \frac{Q}{4\pi\epsilon_0 a} - \frac{Q}{4\pi\epsilon_0 2d} - \frac{p}{4\pi\epsilon_0 4d^2} =$$

$$= \frac{Q}{8\pi\epsilon_0 ad} (2d - a) + \frac{p}{16\pi\epsilon_0 a^2 d^2} (4d^2 - a^2) = \frac{Q}{8\pi\epsilon_0 ad} (2d - a) + \frac{Qa}{64\pi\epsilon_0 d^4} (4d^2 - a^2) \quad \frac{Q}{4\pi\epsilon_0 a} - \frac{Q}{8\pi\epsilon_0 d}$$

$$\left(1 + \frac{a^2}{8d^2}\right)$$

电容值为：

$$C = \frac{Q}{U} = \frac{1}{\frac{1}{8\pi\epsilon_0 ad} (2d - a) + \frac{a}{64\pi\epsilon_0 d^4} (2d - a)(2d + a)} = \frac{8\pi\epsilon_0 ad}{(2d - a) \left[1 + \frac{a^2}{8d^2} (2d + a)\right]}$$

也可以进一步简化为（非必须）：

$$\text{内例 } U = \frac{Q}{4\pi\epsilon_0 a} - \frac{Q}{4\pi\epsilon_0 (2d - a)} + \frac{p}{4\pi\epsilon_0 a^2} - \frac{p}{4\pi\epsilon_0 4d^2}$$

$$= \frac{Q}{4\pi\epsilon_0 a} - \frac{Q}{8\pi\epsilon_0 d} \left[1 - \frac{a^2}{4d^2} + \frac{a^2}{4d^2}\right]$$

$$\text{点电荷: 球心, } U = \frac{Q}{4\pi\epsilon_0 a} - \frac{Q}{8\pi\epsilon_0 d} \left(1 + \frac{a^2}{8d^2}\right)$$

$$\text{内例: } U = \frac{Q}{4\pi\epsilon_0 a} - \frac{Q}{8\pi\epsilon_0 d} \left(1 + \frac{a^2}{8d^2}\right)$$

$$\Rightarrow C = 4\pi\epsilon_0 a \left[1 + \frac{a}{2d} + \frac{a^2}{4d^2} + \frac{a^3}{8d^3} + \frac{a^4}{8d^4}\right]$$

$$\begin{aligned}
 C &= \frac{8\pi\epsilon_0 ad}{(2d-a)} \left[ 1 - \frac{a^2}{8d^3} (2d+a) \right] = 4\pi\epsilon_0 a \left( 1 + \frac{a}{2d} \right) \left[ 1 - \frac{a^2}{8d^3} (2d+a) \right] \\
 &= 4\pi\epsilon_0 a + 2\pi\epsilon_0 \frac{a^2}{d} - \pi\epsilon_0 \frac{a^3}{d^2} - \frac{1}{2} \pi\epsilon_0 \frac{a^4}{d^3} - \frac{1}{2} \pi\epsilon_0 \frac{a^4}{d^3} - \frac{1}{4} \pi\epsilon_0 \frac{a^5}{d^4} \\
 &= 4\pi\epsilon_0 a + 2\pi\epsilon_0 \frac{a^2}{d} - \pi\epsilon_0 \frac{a^3}{d^2} - \pi\epsilon_0 \frac{a^4}{d^3} - \frac{1}{4} \pi\epsilon_0 \frac{a^5}{d^4}
 \end{aligned}$$

4 (4) 探头作用力为四项：  
点电荷之间作用力：

$$F_0 = -\frac{1}{4\pi\epsilon_0} \frac{Q^2}{4d^2} = -\frac{1}{16\pi\epsilon_0} \frac{Q^2}{d^2}$$

点电荷与电偶极子之间作用力(2项)：右边电偶极子受到左边点电荷作用力，右边点电荷受到左边电偶极子的作用力用梯度力计算：

$$F_1 = 2(\vec{p} \cdot \nabla) \vec{E}'_q = 2 \left( p \frac{\partial}{\partial x} \frac{1}{4\pi\epsilon_0} \frac{Q}{x^2} \right)_{x=2d} = -\frac{p}{\pi\epsilon_0} \frac{Q}{(2d)^3} = -\frac{1}{32\pi\epsilon_0} \frac{Q^2 a^3}{d^5}$$

电偶极子与电偶极子之间作用力用梯度力计算：

$$\begin{aligned}
 F_2 &= (\vec{p} \cdot \nabla) \vec{E}'_p = \left( p \frac{\partial}{\partial x} \frac{1}{2\pi\epsilon_0} \frac{p'}{x^3} \right)_{x=2d} = \left( p \frac{\partial}{\partial x} \frac{1}{2\pi\epsilon_0} \frac{Qa^3}{x^5} \right)_{x=2d} \\
 &= -p \frac{5}{2\pi\epsilon_0} \frac{Qa^3}{(2d)^6} = -\frac{5}{2\pi\epsilon_0} \frac{Q^2 a^6}{(2d)^8} = -\frac{5}{512\pi\epsilon_0} \frac{Q^2 a^6}{d^8}
 \end{aligned}$$

合力为：

$$F = -\frac{1}{16\pi\epsilon_0} \frac{Q^2}{d^2} - \frac{1}{32\pi\epsilon_0} \frac{Q^2 a^3}{d^5} - \frac{5}{512\pi\epsilon_0} \frac{Q^2 a^6}{d^8} = -\frac{1}{16\pi\epsilon_0} \frac{Q^2}{d^2} \left( 1 + \frac{a^3}{2d^3} + \frac{5}{32} \frac{a^6}{d^6} \right)$$

4 (5) 初态静电能为(静电能的多极子展开)：

$$\begin{aligned}
 W_0 &= \frac{1}{2} QU - \frac{1}{2} \vec{p} \cdot \vec{E}' = \frac{Q^2}{16\pi\epsilon_0 ad} (2d-a) + \frac{Q^2 a}{128\pi\epsilon_0 d^4} (4d^2 - a^2) - \frac{1}{2} \frac{p^2}{2\pi\epsilon_0 (2d)^3} \\
 &= \frac{Q^2}{16\pi\epsilon_0 ad} (2d-a) + \frac{Q^2 a}{128\pi\epsilon_0 d^4} (4d^2 - a^2) - \frac{Q^2 a^6}{256\pi\epsilon_0 d^7}
 \end{aligned}$$

末态静电能为导体球的自能，即：

$$W = \frac{Q^2}{8\pi\epsilon_0 a}$$

外力做功为：

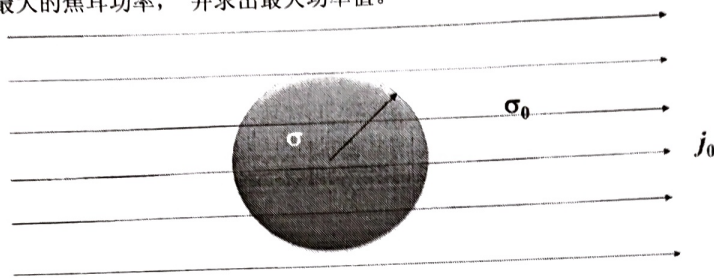
$$A = W - W_0 = \frac{Q^2}{8\pi\epsilon_0 a} - \frac{Q^2}{16\pi\epsilon_0 ad} (2d-a) - \frac{Q^2 a}{128\pi\epsilon_0 d^4} (4d^2 - a^2) + \frac{Q^2 a^6}{256\pi\epsilon_0 d^7}$$

或者简化为(非必须)：

$$A = \frac{Q^2}{16\pi\epsilon_0 d} - \frac{Q^2 a}{128\pi\epsilon_0 d^4} (4d^2 - a^2) + \frac{Q^2 a^6}{256\pi\epsilon_0 d^7} \approx \frac{Q^2}{16\pi\epsilon_0 d} - \frac{Q^2 a}{32\pi\epsilon_0 d^2} = \frac{Q^2}{16\pi\epsilon_0 d} \left( 1 - \frac{a}{2d} \right)$$

$$\frac{Q^2}{16\pi\epsilon_0 d} \left( 1 + \frac{1}{8} \frac{a^3}{d^3} \right)$$

3 (25分) 一个半径为  $R$ 、电导率为  $\sigma$  的导电球放置在电导率为  $\sigma_0$  的无限大空间中，该无限大空间有一个电流密度  $j_0$  均匀分布的电流场，假设球面的电荷在球内产生的电场是均匀电场，在球外产生的电场等效于在球心处的电偶极子产生的电场。求：(1) 球内外的电势分布；(2) 球内外的电场强度分布；(3) 等效电偶极矩和球面电荷分布；(4) 当电导率  $\sigma$  为多少时导体球有最大的焦耳功率，并求出最大功率值。



7 【解】(1) 设球内外的电势分布为：

$$\begin{cases} \varphi_1 = ar \cos \theta \\ \varphi_2 = -E_0 r \cos \theta + \frac{b \cos \theta}{r^2} \end{cases}$$

稳定时，球面处的电场和电流满足边值关系：

$$\begin{cases} E_{1t} = E_{2t} \\ j_{1n} = j_{2n} \end{cases} \quad \begin{cases} E_{1t} = E_{2t} \\ \sigma E_{1n} = \sigma_0 E_{2n} \end{cases}$$

由电势表达式得到球内外电场强度为：

$$\begin{cases} \vec{E}_1 = -\frac{\partial \varphi_1}{\partial r} \vec{e}_r - \frac{1}{r} \frac{\partial \varphi_1}{\partial \theta} \vec{e}_\theta = -a \cos \theta \vec{e}_r + a \sin \theta \vec{e}_\theta \\ \vec{E}_2 = -\frac{\partial \varphi_2}{\partial r} \vec{e}_r - \frac{1}{r} \frac{\partial \varphi_2}{\partial \theta} \vec{e}_\theta = \left( E_0 \cos \theta + \frac{2b \cos \theta}{r^3} \right) \vec{e}_r + \left( -E_0 \sin \theta + \frac{b \sin \theta}{r^3} \right) \vec{e}_\theta \end{cases}$$

代入到球面  $r=R$  处的边值关系中，有：

$$\begin{cases} a \sin \theta = -E_0 \sin \theta + \frac{b \sin \theta}{R^3} \\ -\sigma a \cos \theta = \sigma_0 \left( E_0 \cos \theta + \frac{2b \cos \theta}{R^3} \right) \end{cases}$$

整理一下：

$$\begin{cases} a = -E_0 + \frac{b}{R^3} \\ a = -\frac{\sigma_0}{\sigma} \left( E_0 + \frac{2b}{R^3} \right) \end{cases}$$

$$-E_0 + \frac{b}{R^3} = -\frac{\sigma_0}{\sigma} \left( E_0 + \frac{2b}{R^3} \right)$$

解得：

$$= \frac{2\sigma + \sigma_0}{\sigma + 2\sigma_0} \frac{2\varepsilon_0 j_0}{\sigma_0} \cos\theta - \frac{3\varepsilon_0}{\sigma + 2\sigma_0} j_0 \cos\theta = \left( \frac{4\sigma + 2\sigma_0}{\sigma_0} - 3 \right) \frac{\varepsilon_0}{\sigma + 2\sigma_0} j_0 \cos\theta$$

$$= \frac{(4\sigma - \sigma_0) \varepsilon_0 j_0}{\sigma + 2\sigma_0 \sigma_0} \cos\theta$$

61 (4) 由于球内电场强度是均匀的, 因此焦耳功率密度为:

$$p_{\text{功}} = \sigma E_1^2 = \sigma \left( \frac{3}{\sigma + 2\sigma_0} \right)^2 j_0^2$$

球内总功率为:

$$P_{\text{功}} = p_{\text{功}} \frac{4}{3} \pi R^3 = \sigma \left( \frac{3}{\sigma + 2\sigma_0} \right)^2 j_0^2 \frac{4}{3} \pi R^3$$

为了使该球的功率取极值, 必须:

$$\frac{dP_{\text{功}}}{d\sigma} = \frac{(\sigma + 2\sigma_0)^2 - 2\sigma(\sigma + 2\sigma_0)}{(\sigma + 2\sigma_0)^4} j_0^2 12\pi R^3 = \frac{2\sigma_0 - \sigma}{(\sigma + 2\sigma_0)^3} j_0^2 12\pi R^3 = 0$$

解得:

$$\sigma = 2\sigma_0$$

$$\frac{d^2 P_{\text{功}}}{d\sigma^2} = \frac{-(\sigma + 2\sigma_0)^3 - 3(2\sigma_0 - \sigma)(\sigma + 2\sigma_0)^2}{(\sigma + 2\sigma_0)^6} j_0^2 12\pi R^3$$

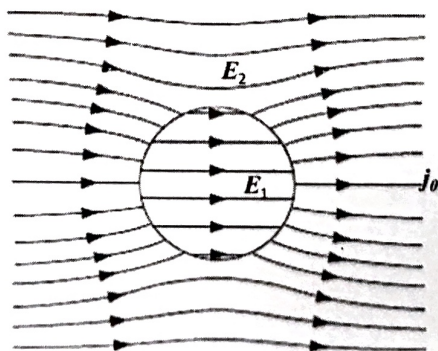
$$= \frac{-\sigma - 2\sigma_0 - 6\sigma_0 + 3\sigma}{(\sigma + 2\sigma_0)^4} j_0^2 12\pi R^3 = \frac{2\sigma - 8\sigma_0}{(\sigma + 2\sigma_0)^4} j_0^2 12\pi R^3$$

$$\left. \frac{d^2 P_{\text{功}}}{d\sigma^2} \right|_{\sigma=2\sigma_0} = \frac{4\sigma_0 - 8\sigma_0}{(2\sigma_0 + 2\sigma_0)^4} j_0^2 12\pi R^3 = -\frac{3}{4\sigma_0^3} j_0^2 \pi R^3 < 0$$

所以  $\sigma = 2\sigma_0$ , 球的功率取极大值。

最大功率为:

$$P_{\text{功}} = 2\sigma_0 \left( \frac{3}{4\sigma_0} \right)^2 j_0^2 \frac{4}{3} \pi R^3 = \frac{3}{2\sigma_0} j_0^2 \pi R^3$$





$$b = \frac{\sigma - \sigma_0}{\sigma + 2\sigma_0} E_0 R^3$$

同理有:

$$a = -E_0 + \frac{b}{R^3} = -E_0 + \frac{(\sigma - \sigma_0)}{\sigma + 2\sigma_0} E_0 = -\frac{3\sigma_0}{\sigma + 2\sigma_0} E_0$$

代入到电势中, 由于  $\vec{E}_0 = \vec{j}_0 / \sigma_0$ , 得到:

$$\begin{cases} \varphi_1 = -\frac{3\sigma_0}{\sigma + 2\sigma_0} E_0 r \cos\theta = -\frac{3}{\sigma + 2\sigma_0} j_0 r \cos\theta \\ \varphi_2 = -\frac{j_0}{\sigma_0} r \cos\theta + \frac{\sigma - \sigma_0}{\sigma + 2\sigma_0} \frac{R^3}{r^2} \frac{j_0}{\sigma_0} \cos\theta \end{cases}$$

7' (2) 球内电场为:

$$\begin{aligned} \vec{E}_1 &= \frac{3\sigma_0}{\sigma + 2\sigma_0} E_0 \cos\theta \vec{e}_r - \frac{3\sigma_0}{\sigma + 2\sigma_0} E_0 \sin\theta \vec{e}_\theta \\ &= \frac{3\sigma_0}{\sigma + 2\sigma_0} (E_0 \cos\theta \vec{e}_r - E_0 \sin\theta \vec{e}_\theta) = \frac{3\sigma_0}{\sigma + 2\sigma_0} \vec{E}_0 = \frac{3}{\sigma + 2\sigma_0} \vec{j}_0 \end{aligned}$$

球外电场强度为:

$$\begin{aligned} \vec{E}_2 &= \left(1 + \frac{\sigma - \sigma_0}{\sigma + 2\sigma_0} \frac{2R^3}{r^3}\right) E_0 \cos\theta \vec{e}_r + \left(\frac{\sigma - \sigma_0}{\sigma + 2\sigma_0} \frac{R^3}{r^3} - 1\right) E_0 \sin\theta \vec{e}_\theta \\ &= \left(1 + \frac{\sigma - \sigma_0}{\sigma + 2\sigma_0} \frac{2R^3}{r^3}\right) \frac{j_0}{\sigma_0} \cos\theta \vec{e}_r + \left(\frac{\sigma - \sigma_0}{\sigma + 2\sigma_0} \frac{R^3}{r^3} - 1\right) \frac{j_0}{\sigma_0} \sin\theta \vec{e}_\theta \end{aligned}$$

5' (3) 求等效电偶极矩, 只需比较球外电势的第二项与标准的电偶极矩电势表达式, 即:

$$\frac{\sigma - \sigma_0}{\sigma + 2\sigma_0} \frac{R^3}{r^2} \frac{j_0}{\sigma_0} \cos\theta = \frac{1}{4\pi\epsilon_0} \frac{p \cos\theta}{r^2}$$

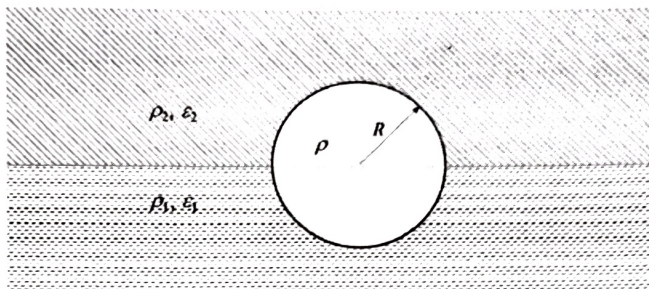
得到:

$$3' \quad p = \frac{\sigma - \sigma_0}{\sigma + 2\sigma_0} 4\pi\epsilon_0 R^3 \frac{j_0}{\sigma_0}$$

球面的电荷密度为:

$$\begin{aligned} 2' \quad \sigma_c &= \epsilon_0 (E_{2n} - E_{1n}) = \left(1 + \frac{\sigma - \sigma_0}{\sigma + 2\sigma_0} \frac{2R^3}{r^3}\right) \frac{\epsilon_0 j_0}{\sigma_0} \cos\theta - \frac{3\epsilon_0}{\sigma + 2\sigma_0} j_0 \cos\theta \\ &= \left(1 + \frac{2\sigma - 2\sigma_0}{\sigma + 2\sigma_0} - \frac{3\sigma_0}{\sigma + 2\sigma_0}\right) \frac{\epsilon_0 j_0}{\sigma_0} \cos\theta \\ &= \frac{\sigma + 2\sigma_0 + 2\sigma - 2\sigma_0 - 3\sigma_0}{\sigma + 2\sigma_0} \frac{\epsilon_0 j_0}{\sigma_0} \cos\theta \\ &= \frac{3\sigma - 3\sigma_0}{\sigma + 2\sigma_0} \frac{\epsilon_0 j_0}{\sigma_0} \cos\theta = \frac{\sigma - \sigma_0}{\sigma + 2\sigma_0} 3 \frac{\epsilon_0 j_0}{\sigma_0} \cos\theta. \end{aligned}$$

4. (25分) 一个半径为 $R$ , 质量密度为 $\rho$ 的导体球的一半悬浮在密度为 $\rho_1$ , ( $\rho_1 > 2\rho$ ) 的液体中, 液体的相对介电常数为 $\epsilon_1$ , 球的上半部分是另一种液体, 质量密度为 $\rho_2$  ( $\rho_2 < \rho$ ), 相对介电常数为 $\epsilon_2$ , 且 $\epsilon_2 < \epsilon_1$ . (1) 给导体球带电量 $Q$ 为多少时, 系统达到平衡; (2) 平衡时上半球面和下半球面的自由电荷量、极化电荷的量和总电荷量; (3) 系统的电容值; (4) 若介质均为导电介质, 电导率分别为 $\sigma_{1\text{电}}$ 和 $\sigma_{2\text{电}}$ , 则球面上自由电荷 $Q$ 最终会到无限远处(设漏电过程中球的位置不变), 则该漏电过程中消耗的总能量为多少? (5) 漏电电阻为多少? 漏电时间常数 $\tau$ 为多少?



【解】(1) 介质分界面为垂直于等势面情况, 两种介质中的电场强度相同, 设导体球带电量为 $Q$ , 利用高斯定理, 有:

$$2\pi r^2 D_1 + 2\pi r^2 D_2 = Q$$

因为:  $D_1 = \epsilon_0 \epsilon_1 E_1$ ,  $D_2 = \epsilon_0 \epsilon_2 E_2$ , 且  $E_1 = E_2 = E$

$$2\pi R^2 \epsilon_0 (\epsilon_1 + \epsilon_2) E = Q$$

$$E = \frac{Q}{2\pi r^2 \epsilon_0 (\epsilon_1 + \epsilon_2)}$$

因此得:

$$D_1 = \epsilon_0 \epsilon_1 E_1 = \frac{\epsilon_1 Q}{2\pi r^2 (\epsilon_1 + \epsilon_2)}, \quad D_2 = \epsilon_0 \epsilon_2 E_2 = \frac{\epsilon_2 Q}{2\pi r^2 (\epsilon_1 + \epsilon_2)}$$

下半球球面的静电压强为:

$$\begin{aligned} p_1 &= \frac{1}{2} \vec{E}_1 \cdot \vec{D}_1 = \frac{1}{2} E_1 D_1 = \frac{1}{2} \cdot \frac{Q}{2\pi r^2 \epsilon_0 (\epsilon_1 + \epsilon_2)} \cdot \frac{\epsilon_1 Q}{2\pi r^2 (\epsilon_1 + \epsilon_2)} \\ &= \frac{\epsilon_1 Q^2}{8\pi^2 \epsilon_0 R^4 (\epsilon_1 + \epsilon_2)^2} \end{aligned}$$

方向沿径向向外。同理, 上半球球面的静电压强为:

$$p_2 = \frac{\epsilon_2 Q^2}{8\pi^2 \epsilon_0 R^4 (\epsilon_1 + \epsilon_2)^2}$$

方向沿径向向外。

上下半球总静电力为:

$$F_E = p_1 \pi R^2 - p_2 \pi R^2 = \frac{(\varepsilon_1 - \varepsilon_2) Q^2}{8\pi \varepsilon_0 R^2 (\varepsilon_1 + \varepsilon_2)^2}$$

方向向下。

导体球在两种液体界面处的浮力为:

$$F_H = \frac{2}{3} \pi R^3 (\rho_1 + \rho_2) g$$

导体球的总重量为:

$$F_G = \frac{4}{3} \pi R^3 \rho g$$

根据平衡条件, 得:

$$\frac{(\varepsilon_1 - \varepsilon_2) Q^2}{8\pi \varepsilon_0 R^2 (\varepsilon_1 + \varepsilon_2)^2} + \frac{4}{3} \pi R^3 \rho g = \frac{2}{3} \pi R^3 (\rho_1 + \rho_2) g$$

$$Q^2 = \frac{16}{3} \varepsilon_0 \pi^2 R^5 \frac{(\varepsilon_1 + \varepsilon_2)^2}{(\varepsilon_1 - \varepsilon_2)} (\rho_1 + \rho_2 - 2\rho) g$$

$$Q = 4\pi R^2 (\varepsilon_1 + \varepsilon_2) \sqrt{\frac{\varepsilon_0 R}{3(\varepsilon_1 - \varepsilon_2)} (\rho_1 + \rho_2 - 2\rho) g}$$

(2) 上下半球面的自由电荷面密度分布为:

$$\sigma_1 = D_1|_{r=R} = \frac{\varepsilon_1 Q}{2\pi R^2 (\varepsilon_1 + \varepsilon_2)} = 2\varepsilon_1 \sqrt{\frac{\varepsilon_0 R}{3(\varepsilon_1 - \varepsilon_2)} (\rho_1 + \rho_2 - 2\rho) g}$$

$$\sigma_2 = D_2|_{r=R} = \frac{\varepsilon_2 Q}{2\pi R^2 (\varepsilon_1 + \varepsilon_2)} = 2\varepsilon_2 \sqrt{\frac{\varepsilon_0 R}{3(\varepsilon_1 - \varepsilon_2)} (\rho_1 + \rho_2 - 2\rho) g}$$

上下半球面的总自由电荷量分别为:

$$Q_1 = \sigma_1 2\pi R^2 = \frac{\varepsilon_1 Q}{(\varepsilon_1 + \varepsilon_2)} = 4\pi R^2 \varepsilon_1 \sqrt{\frac{\varepsilon_0 R}{3(\varepsilon_1 - \varepsilon_2)} (\rho_1 + \rho_2 - 2\rho) g}$$

$$Q_2 = \sigma_2 2\pi R^2 = \frac{\varepsilon_2 Q}{(\varepsilon_1 + \varepsilon_2)} = 4\pi R^2 \varepsilon_2 \sqrt{\frac{\varepsilon_0 R}{3(\varepsilon_1 - \varepsilon_2)} (\rho_1 + \rho_2 - 2\rho) g}$$

导体球上的总自由电荷为 (非必须):

$$Q_1 + Q_2 = \frac{\varepsilon_1 Q}{(\varepsilon_1 + \varepsilon_2)} + \frac{\varepsilon_2 Q}{(\varepsilon_1 + \varepsilon_2)} = Q = 4\pi R^2 (\varepsilon_1 + \varepsilon_2) \sqrt{\frac{\varepsilon_0 R}{3(\varepsilon_1 - \varepsilon_2)} (\rho_1 + \rho_2 - 2\rho) g}$$

(注:  $Q$  值可以代入, 也可以不代入)

两种介质中的极化强度为:

$$P_1 = \varepsilon_0 (\varepsilon_1 - 1) E = \frac{(\varepsilon_1 - 1) Q}{2\pi r^2 \varepsilon_0 (\varepsilon_1 + \varepsilon_2)}, \quad P_2 = \varepsilon_0 (\varepsilon_2 - 1) E = \frac{(\varepsilon_2 - 1) Q}{2\pi r^2 \varepsilon_0 (\varepsilon_1 + \varepsilon_2)}$$

上下半球面的自由电荷面密度分布为:

$$\sigma'_1 = -P_1|_{r=R} = \frac{(\varepsilon_1 - 1)Q}{2\pi R^2(\varepsilon_1 + \varepsilon_2)}, \quad \sigma'_2 = -P_2|_{r=R} = \frac{(\varepsilon_2 - 1)Q}{2\pi R^2(\varepsilon_1 + \varepsilon_2)}$$

上下半球面的极化电荷量分别为:

$$Q'_1 = 2\pi R^2 \sigma'_1 = -\frac{(\varepsilon_1 - 1)Q}{(\varepsilon_1 + \varepsilon_2)}, \quad Q'_2 = 2\pi R^2 \sigma'_2 = -\frac{(\varepsilon_2 - 1)Q}{(\varepsilon_1 + \varepsilon_2)}$$

上下球面总极化电荷量为 (非必须):

$$Q' = Q'_1 + Q'_2 = -\frac{(\varepsilon_1 - 1)Q}{(\varepsilon_1 + \varepsilon_2)} - \frac{(\varepsilon_2 - 1)Q}{(\varepsilon_1 + \varepsilon_2)} = -\frac{(\varepsilon_1 - 1) + (\varepsilon_2 - 1)}{(\varepsilon_1 + \varepsilon_2)} Q = -\frac{(\varepsilon_1 + \varepsilon_2 - 2)}{(\varepsilon_1 + \varepsilon_2)} Q$$

上下半球面的总电荷量为:

$$Q_{\text{上}} = Q_1 + Q'_1 = \frac{\varepsilon_1 Q}{(\varepsilon_1 + \varepsilon_2)} - \frac{(\varepsilon_1 - 1)Q}{(\varepsilon_1 + \varepsilon_2)} = \frac{Q}{(\varepsilon_1 + \varepsilon_2)}$$

$$Q_{\text{下}} = Q_2 + Q'_2 = \frac{\varepsilon_2 Q}{(\varepsilon_1 + \varepsilon_2)} - \frac{(\varepsilon_2 - 1)Q}{(\varepsilon_1 + \varepsilon_2)} = \frac{Q}{(\varepsilon_1 + \varepsilon_2)}$$

式子 $Q$ 为(1)中的电量(可以代入,也可以不代入)。

即球面总电荷或总电荷密度分布处处相同。

(3) 导体球的电势为:

$$U = -\int_{\infty}^R E dr = \frac{Q}{2\pi r \varepsilon_0 (\varepsilon_1 + \varepsilon_2)} \Big|_{\infty}^R = \frac{Q}{2\pi R \varepsilon_0 (\varepsilon_1 + \varepsilon_2)}$$

等效电容为:

$$C = \frac{Q}{U} = 2\pi R \varepsilon_0 (\varepsilon_1 + \varepsilon_2)$$

(4) 解法I: 初始时刻导体球带电产生的静电能为:

$$\begin{aligned} W &= W_1 + W_2 = \frac{1}{2} \varepsilon_0 \varepsilon_1 \iiint_{\text{上}} E^2 dV + \frac{1}{2} \varepsilon_0 \varepsilon_2 \iiint_{\text{下}} E^2 dV \\ &= \frac{1}{2} \varepsilon_0 (\varepsilon_1 + \varepsilon_2) \iiint_{\text{上}} E^2 dV = \frac{1}{2} \varepsilon_0 (\varepsilon_1 + \varepsilon_2) \int_R^{\infty} \left[ \frac{Q}{2\pi r^2 \varepsilon_0 (\varepsilon_1 + \varepsilon_2)} \right]^2 2\pi r^2 dr \\ &= \frac{1}{4\pi \varepsilon_0 (\varepsilon_1 + \varepsilon_2)} \int_R^{\infty} \frac{1}{r^2} dr = \frac{Q^2}{4\pi \varepsilon_0 R (\varepsilon_1 + \varepsilon_2)} = \frac{4}{3} \pi R^4 \frac{(\varepsilon_1 + \varepsilon_2)}{(\varepsilon_2 - \varepsilon_1)} (\rho_1 + \rho_2 - 2\rho) g \end{aligned}$$

这部分静电能在电荷流到无限远时, 变为零, 即全部转化为焦耳热。

解法II: 用电容器储能计算, 即:

$$W = \frac{1}{2} \frac{Q^2}{C} = \frac{Q^2}{4\pi R \varepsilon_0 (\varepsilon_1 + \varepsilon_2)} = \frac{4}{3} \pi R^4 \frac{(\varepsilon_1 + \varepsilon_2)}{(\varepsilon_2 - \varepsilon_1)} (\rho_1 + \rho_2 - 2\rho) g$$

两者算法答案相同。

(5) 上下半球的电容和漏电阻分别为:

$$C_1 = 2\pi R \epsilon_0 \epsilon_1, \quad R_1 C_1 = \frac{\epsilon_0 \epsilon_1}{\sigma_{1\text{导}}}, \quad R_1 = \frac{\epsilon_0 \epsilon_1}{\sigma_{1\text{导}} C_1} = \frac{\epsilon_0 \epsilon_1}{\sigma_{1\text{导}} 2\pi R \epsilon_0 \epsilon_1} = \frac{1}{\sigma_{1\text{导}} 2\pi R}$$

$$C_2 = 2\pi R \epsilon_0 \epsilon_2, \quad R_2 C_2 = \frac{\epsilon_0 \epsilon_2}{\sigma_{2\text{导}}}, \quad R_2 = \frac{\epsilon_0 \epsilon_2}{\sigma_{2\text{导}} C_2} = \frac{\epsilon_0 \epsilon_2}{\sigma_{2\text{导}} 2\pi R \epsilon_0 \epsilon_2} = \frac{1}{\sigma_{2\text{导}} 2\pi R}$$

总电阻为上下两部分电阻并联，即：

$$R_{\text{阻}} = \frac{R_1 R_2}{R_1 + R_2} = \frac{\frac{1}{\sigma_{1\text{导}} 2\pi R} \cdot \frac{1}{\sigma_{2\text{导}} 2\pi R}}{\frac{1}{\sigma_{1\text{导}} 2\pi R} + \frac{1}{\sigma_{2\text{导}} 2\pi R}} = \frac{1}{2\pi R (\sigma_{1\text{导}} + \sigma_{2\text{导}})}$$

漏电时间常数  $\tau$  为：

$$\tau = RC = \frac{1}{2\pi R (\sigma_{1\text{导}} + \sigma_{2\text{导}})} 2\pi R \epsilon_0 (\epsilon_1 + \epsilon_2) = \frac{\epsilon_0 (\epsilon_1 + \epsilon_2)}{(\sigma_{1\text{导}} + \sigma_{2\text{导}})}$$