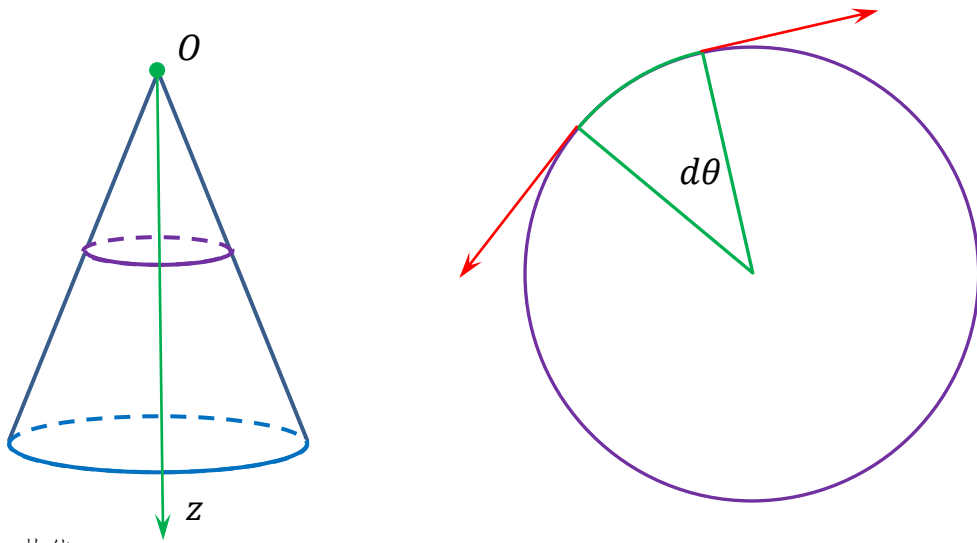


一、 (20%)



解法 1: 势能

$$V = -mgz + \frac{1}{2}k(2\pi z \tan \alpha - l)^2$$

$$\delta V = \{-mg + k(2\pi z \tan \alpha - l)2\pi \tan \alpha\} \delta z = 0$$

$$2\pi z \tan \alpha = \frac{mg}{2\pi k \tan \alpha} + l$$

解法 2: 如图, 取一小段橡皮绳微元来分析。这段绳受到三个主动力的作用。其中有两个主动力为弹性张力, 其方向在水平面内, 大小为  $k(2\pi z \tan \alpha - l)$ ; 这两个力的合力为  $2 \times k(2\pi z \tan \alpha - l) \cos(\frac{\pi}{2} - \frac{d\theta}{2}) = k(2\pi z \tan \alpha - l)d\theta$ , 方向在水平面内, 指向橡皮圈的圆心。另一个主动力是重力, 垂直向下, 大小为  $\frac{m}{2\pi}d\theta g$ 。

设橡皮绳微元沿侧面向下发生虚位移, 其  $z$  分量为  $\delta z$ , 水平分量为  $\delta z \tan \alpha$ , 由虚功原理

$$\sum_i \vec{F}_i \cdot \delta \vec{r}_i = 0$$

$$-k(2\pi z \tan \alpha - l)d\theta \cdot \delta z \tan \alpha + \frac{m}{2\pi}d\theta g \cdot \delta z = 0$$

$$-k(2\pi z \tan \alpha - l) \tan \alpha + \frac{mg}{2\pi} = 0$$

$$2\pi z \tan \alpha = \frac{mg}{2\pi k \tan \alpha} + l$$

二、 (20 分)

解:

(1) 变换为

$$t' = t$$

$$\vec{r}'_i = \vec{r}_i + \vec{v}_0 t$$

$$\dot{\vec{r}}'_i = \dot{\vec{r}}_i + \vec{v}_0$$

得

$$L' dt' = \left\{ \frac{1}{2} m_1 \dot{\vec{r}}_1'^2 + \frac{1}{2} m_2 \dot{\vec{r}}_2'^2 - V(\vec{r}'_1 - \vec{r}'_2) \right\} dt'$$

$$= \left\{ \frac{1}{2} m_1 (\dot{\vec{r}}_1 + \vec{\epsilon})^2 + \frac{1}{2} m_2 (\dot{\vec{r}}_2 + \vec{\epsilon})^2 - V(\vec{r}_1 - \vec{r}_2) \right\} dt$$

$$= L dt + (m_1 \dot{\vec{r}}_1 + m_2 \dot{\vec{r}}_2) \cdot \vec{\epsilon} dt + \mathcal{O}(\epsilon^2)$$

$$= L dt + d\{(m_1 \vec{r}_1 + m_2 \vec{r}_2) \cdot \vec{\epsilon}\} + \mathcal{O}(\epsilon^2)$$

所以是准对称变换。

(2) 由

$$\Delta t = 0, \quad \Delta \vec{r}_i = \vec{\epsilon} t, \quad \Delta \varphi = (m_1 \vec{r}_1 + m_2 \vec{r}_2) \cdot \vec{\epsilon}$$

$$\vec{p}_i = \frac{\partial L}{\partial \dot{\vec{r}}_i} = m_i \dot{\vec{r}}_i (i \text{ 不求和})$$

$$-H\Delta t + \vec{p}_i \cdot \Delta \vec{r}_i - \Delta \varphi = (m_1 \dot{\vec{r}}_1 + m_2 \dot{\vec{r}}_2) \cdot \vec{\epsilon} t - (m_1 \vec{r}_1 + m_2 \vec{r}_2) \cdot \vec{\epsilon}$$

守恒量为

$$(m_1 \dot{\vec{r}}_1 + m_2 \dot{\vec{r}}_2)t - (m_1 \vec{r}_1 + m_2 \vec{r}_2) = M(\vec{v}_c t - \vec{r}_c) = M\vec{r}_c(0)$$

三、 (20 分)

解：光程

$$S[r, \theta] = \int_{z_1}^{z_2} n_0 \sqrt{1 - \alpha^2 r} \sqrt{1 + r'^2 + r^2 \theta'^2} dz$$

拉氏量为

$$L = n_0 \sqrt{(1 - \alpha^2 r)(1 + r'^2 + r^2 \theta'^2)}$$

拉氏方程

$$\frac{d}{dz} \frac{r' \sqrt{1 - \alpha^2 r}}{\sqrt{1 + r'^2 + r^2 \theta'^2}} + \frac{\alpha^2(1 + r'^2) + (3\alpha^2 r^2 - 2r)\theta'^2}{2\sqrt{(1 - \alpha^2 r)(1 + r'^2 + r^2 \theta'^2)}} = 0$$

$$\frac{d}{dz} \left( r^2 \theta' \sqrt{\frac{1 - \alpha^2 r}{1 + r'^2 + r^2 \theta'^2}} \right) = 0$$

四、 (20 分)

解：莫培督原理

$$0 = \delta \int_{s_1}^{s_2} ds = \delta \int_{\theta_1}^{\theta_2} \sqrt{2m[E - V(r)]} \sqrt{r'^2 + r^2} d\theta$$

$$= \delta \int_{\theta_1}^{\theta_2} \sqrt{2m(E - \alpha u)} \sqrt{\frac{1}{u^4} u'^2 + \frac{1}{u^2}} d\theta$$

$$= \delta \int_{\theta_1}^{\theta_2} \frac{1}{u^2} \sqrt{2m(E - \alpha u)} \sqrt{u'^2 + u^2} d\theta$$

$$\frac{d}{d\theta} \frac{1}{u^2} \sqrt{E - \alpha u} \frac{u'}{\sqrt{u'^2 + u^2}} + \frac{2}{u^3} \sqrt{E - \alpha u} \sqrt{u'^2 + u^2} + \frac{1}{2u^2} \frac{\alpha}{\sqrt{E - \alpha u}} \sqrt{u'^2 + u^2} - \frac{1}{u} \frac{1}{\sqrt{E - \alpha u}} \frac{1}{\sqrt{u'^2 + u^2}} = 0$$

$$\frac{2}{u} \frac{\sqrt{E - \alpha u}}{\sqrt{u'^2 + u^2}} + \frac{\alpha}{2\sqrt{E - \alpha u} \sqrt{u'^2 + u^2}} + \frac{1}{u^2} \sqrt{E - \alpha u} \left( \frac{d}{d\theta} \frac{u'}{\sqrt{u'^2 + u^2}} - \frac{u}{\sqrt{u'^2 + u^2}} \right) = 0$$

$$\frac{2}{u} \frac{\sqrt{E - \alpha u}}{\sqrt{u'^2 + u^2}} + \frac{\alpha}{2\sqrt{E - \alpha u} \sqrt{u'^2 + u^2}} + \frac{\sqrt{E - \alpha u}}{u^2} \left( \frac{u''}{\sqrt{u'^2 + u^2}} - \frac{u'(u'' + u)}{(u'^2 + u^2)^{\frac{3}{2}}} - \frac{u}{\sqrt{u'^2 + u^2}} \right) = 0$$

$$\frac{2}{u} \frac{\sqrt{E - \alpha u}}{\sqrt{u'^2 + u^2}} + \frac{\alpha}{2\sqrt{E - \alpha u} \sqrt{u'^2 + u^2}} + \frac{\sqrt{E - \alpha u}}{u^2} \frac{1}{(u'^2 + u^2)^{\frac{3}{2}}} (uu'' - 2u'^2 - u^2) = 0$$

$$4(E - \alpha u)(u'^2 + u^2) + \alpha u(u'^2 + u^2) + 2(E - \alpha u)(uu'' - 2u'^2 - u^2) = 0$$

$$2(E - \alpha u)(u'^2 + u^2) + \alpha u(u'^2 + u^2) + 2(E - \alpha u)(uu'' - u'^2) = 0$$

$$2(E - \alpha u)uu'' + \alpha uu'^2 + 2Eu^2 - \alpha u^3 = 0$$

$$2(E - \alpha u)u'' + \alpha(u'^2 + u^2) + 2(E - \alpha u)u = 0$$

注：如果把广义能量积分

$$J = \sqrt{2m(E - \alpha u)} \frac{1}{\sqrt{u'^2 + u^2}}$$

$$\frac{u'^2 + u^2}{2(E - \alpha u)} = \frac{m}{J^2}$$

代入，则方程化简为 Binet 方程

$$u'' + u = \frac{m\alpha}{J^2}$$

五、 (20%)

解：

(1)

$$\begin{aligned}\partial_t \psi &= \dot{q}_1(t)(x^2 - a^2) + \dot{q}_2(t)(x^2 - a^2)x \\ \partial_x \psi &= q_1(t) \cdot 2x + q_2(t)(3x^2 - a^2)\end{aligned}$$

作用量

$$S = \int_{t_1}^{t_2} dt \int_{-a}^a dx \mathcal{L}$$

$$\begin{aligned}L &= \int_{-a}^a dx \mathcal{L} \\ &= \frac{1}{2} \rho \dot{q}_1^2 \int_{-a}^a (x^2 - a^2)^2 dx + \rho \dot{q}_1 \dot{q}_2 \int_{-a}^a (x^2 - a^2)^2 x dx + \frac{1}{2} \rho \dot{q}_2^2 \int_{-a}^a (x^2 - a^2)^2 x^2 dx \\ &\quad - \frac{1}{2} Y q_1^2 \int_{-a}^a 4x^2 dx - Y q_1 q_2 \int_{-a}^a 2x(3x^2 - a^2) dx - \frac{1}{2} Y q_2^2 \int_{-a}^a (3x^2 - a^2)^2 dx \\ &= \frac{1}{2} \rho \dot{q}_1^2 \cdot \frac{16a^5}{15} + \rho \dot{q}_1 \dot{q}_2 \cdot 0 + \frac{1}{2} \rho \dot{q}_2^2 \cdot \frac{16a^7}{105} - \frac{1}{2} Y q_1^2 \cdot \frac{8a^3}{3} - Y q_1 q_2 \cdot 0 - \frac{1}{2} Y q_2^2 \cdot \frac{8a^5}{5} \\ &= \frac{8\rho a^5}{105} (7\dot{q}_1^2 + a^2 \dot{q}_2^2) - \frac{4Ya^3}{15} (5q_1^2 + 3a^2 q_2^2)\end{aligned}$$

(2) 运动方程

$$\begin{aligned}2\rho a^2 \ddot{q}_1 + 5Y q_1 &= 0 \\ 2\rho a^2 \ddot{q}_2 + 21Y q_2 &= 0\end{aligned}$$