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1. Define the density operator

$$\hat{\rho}(t) \equiv \sum_j P_j |\Psi_j(t)\rangle \langle \Psi_j(t)|, \quad (1)$$

where P_j is the probability of finding the ensemble being in the state $|\Psi_j(t)\rangle$ at time t . Note here $\{P_j\}$ is the classical probability with mixing quantum states $\{|\Psi_j(t)\rangle\}$ in an incoherent way. Show that the density operator satisfies the Liouville–von Neumann equation,

$$i\hbar \frac{d}{dt} \hat{\rho}(t) = [\hat{H}, \hat{\rho}(t)], \quad (2)$$

with the given Hamiltonian \hat{H} . Furthermore, show that density operators at different times are related by a unitary transformation. Find the unitary operator.

2. The von Neumann entropy is defined via

$$S_{\text{vN}} \equiv -k_B \text{Tr}(\hat{\rho} \ln \hat{\rho}), \quad (3)$$

with k_B being the Boltzmann constant. For an isolate system, when it approaches to the thermodynamic equilibrium, the entropy reaches maximum, that is

$$\delta S = 0, \quad (4)$$

and the density operator satisfies

$$\dot{\hat{\rho}} = -\frac{i}{\hbar} [\hat{H}, \hat{\rho}] = 0. \quad (5)$$

Prove that for this system,

$$\hat{\rho} = \frac{1}{\Omega} \sum_j |\psi_j\rangle \langle \psi_j|, \quad (6)$$

with $|\psi_j\rangle$ ($j = 1, 2, \dots$) being the eigenstates of the system Hamiltonian and Ω the total number of states. (*Hint: The condition $\text{Tr} \hat{\rho} = 1$ may be helpful.*)