

### 一、15%

解：以角度为广义坐标。平衡时，小球到圆心连线之间的夹角为120°。

设微振动坐标为 $\theta_1, \theta_2, \theta_3$ 。

系统的动能

$$T = \frac{1}{2}ma^2(\dot{\theta}_1^2 + \dot{\theta}_2^2 + \dot{\theta}_3^2)$$

势能

$$\begin{aligned} V &= \frac{1}{2}ka^2\{(\theta_1 - \theta_2)^2 + (\theta_2 - \theta_3)^2 + (\theta_3 - \theta_1)^2\} \\ &= \frac{1}{2}ka^2\{2\theta_1^2 + 2\theta_2^2 + 2\theta_3^2 - 2\theta_1\theta_2 - 2\theta_1\theta_3 - 2\theta_2\theta_3\} \end{aligned}$$

惯性矩阵

$$M = ma^2 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

3分；只写 T, 未能写出 M 的 2 分

刚度矩阵

$$K = ka^2 \begin{pmatrix} 2 & -1 & -1 \\ -1 & 2 & -1 \\ -1 & -1 & 2 \end{pmatrix}$$

3分；只写 V, 未能写出 K 的 2 分

广义特征方程

$$K\vec{x} = \omega^2 M \vec{x}$$

3分

广义特征值

$$\det(K - \omega^2 M) = 0 \Rightarrow \omega^2 = 0, \frac{3k}{m}$$

3分

广义特征矢

$$(1,1,1)^T, \quad (1,-1,0)^T, \quad (1,0,-1)^T$$

施密特正角化，( $\vec{x}_2$ 和 $\vec{x}_3$ 不唯一)

$$\vec{x}_1 = \frac{1}{\sqrt{3}ma}(1,1,1)^T, \quad \vec{x}_2 = \frac{1}{\sqrt{2}ma}(1,-1,0)^T, \quad \vec{x}_3 = \frac{1}{\sqrt{6}ma}(1,1,-2)^T$$

模态矩阵

$$A = (\vec{x}_1, \vec{x}_2, \vec{x}_3) = \frac{1}{\sqrt{6}ma} \begin{pmatrix} \sqrt{2} & \sqrt{3} & 1 \\ \sqrt{2} & -\sqrt{3} & 1 \\ \sqrt{2} & 0 & -2 \end{pmatrix}$$

逆矩阵

$$A^{-1} = A^T M = \frac{\sqrt{ma}}{\sqrt{6}} \begin{pmatrix} \sqrt{2} & \sqrt{2} & \sqrt{2} \\ \sqrt{3} & -\sqrt{3} & 0 \\ 1 & 1 & -2 \end{pmatrix}$$

简正坐标

$$\vec{\xi} = A^{-1} \vec{\eta} = \frac{\sqrt{ma}}{\sqrt{6}} \begin{pmatrix} \sqrt{2} & \sqrt{2} & \sqrt{2} \\ \sqrt{3} & -\sqrt{3} & 0 \\ 1 & 1 & -2 \end{pmatrix} \begin{pmatrix} \theta_1 \\ \theta_2 \\ \theta_3 \end{pmatrix} = \begin{pmatrix} \frac{\sqrt{ma}}{\sqrt{3}}(\theta_1 + \theta_2 + \theta_3) \\ \frac{\sqrt{ma}}{\sqrt{2}}(\theta_1 - \theta_2) \\ \frac{\sqrt{ma}}{\sqrt{6}}(\theta_1 + \theta_2 - \theta_3) \end{pmatrix}$$

简正坐标归一化常数可自由约定；第二、第三个简正坐标的表达式不唯一。

简正模式 3 分；为正角化得 1 分

二、20%

解：(1) 费马原理的拉氏函数是

$$L = n \frac{ds}{dz} = n_0 \sqrt{(1 - \alpha^2 r^2)(r'^2 + r^2 \theta'^2 + 1)}$$

4分

(2) 广义动量  $p_\theta$  守恒，

$$p_\theta = \frac{\partial L}{\partial \theta'} = n_0 \sqrt{\frac{1 - \alpha^2 r^2}{r'^2 + r^2 \theta'^2 + 1}} r^2 \theta'$$

2分

广义能量守恒，

$$\begin{aligned} p_r &= n_0 \sqrt{\frac{1 - \alpha^2 r^2}{r'^2 + r^2 \theta'^2 + 1}} r' \\ \Rightarrow E &= p_r r' + p_\theta \theta' - L = -n_0 \sqrt{\frac{1 - \alpha^2 r^2}{r'^2 + r^2 \theta'^2 + 1}} \end{aligned}$$

2分

本小题共 4 分

(3) 等效拉氏函数是

$$L_{\text{eff}} = L - p_\theta \theta' = n_0 \sqrt{1 - \alpha^2 r^2} \frac{r'^2 + 1}{\sqrt{r'^2 + r^2 \theta'^2 + 1}}$$

2分

消去  $r^2 \theta'^2$ ，

$$\begin{aligned} p_\theta &= n_0 \sqrt{\frac{1 - \alpha^2 r^2}{r'^2 + r^2 \theta'^2 + 1}} r^2 \theta' \Rightarrow \frac{r^2 \theta'^2}{r'^2 + r^2 \theta'^2 + 1} = \frac{p_\theta^2}{n_0^2 (1 - \alpha^2 r^2) r^2} \\ \Rightarrow \sqrt{\frac{r'^2 + 1}{r'^2 + r^2 \theta'^2 + 1}} &= \sqrt{\frac{n_0^2 (1 - \alpha^2 r^2) r^2 - p_\theta^2}{n_0^2 (1 - \alpha^2 r^2) r^2}} \\ \Rightarrow L_{\text{eff}} &= \sqrt{n_0^2 (1 - \alpha^2 r^2) - \frac{p_\theta^2}{r^2} \sqrt{r'^2 + 1}} \end{aligned}$$

2分

本小题共 4 分

(4) 前面已得

$$H = p_r r' + p_\theta \theta' - L = -n_0 \sqrt{\frac{1 - \alpha^2 r^2}{r'^2 + r^2 \theta'^2 + 1}}$$

2分

利用前面的结论作变量替换，

$$\begin{cases} \frac{r^2 \theta'^2}{r'^2 + r^2 \theta'^2 + 1} = \frac{p_\theta^2}{n_0^2 (1 - \alpha^2 r^2) r^2} \Rightarrow \frac{r'^2 + r^2 \theta'^2}{r'^2 + r^2 \theta'^2 + 1} = \frac{1}{n_0^2 (1 - \alpha^2 r^2)} \left( p_r^2 + \frac{p_\theta^2}{r^2} \right) \\ \frac{r'^2}{r'^2 + r^2 \theta'^2 + 1} = \frac{p_r^2}{n_0^2 (1 - \alpha^2 r^2)} \\ \Rightarrow \frac{1}{r'^2 + r^2 \theta'^2 + 1} = \frac{1}{n_0^2 (1 - \alpha^2 r^2)} \left( n_0^2 (1 - \alpha^2 r^2) - p_r^2 - \frac{p_\theta^2}{r^2} \right) \end{cases}$$

$$H = - \sqrt{n_0^2 (1 - \alpha^2 r^2) - p_r^2 - \frac{p_\theta^2}{r^2}}$$

2分

正则方程

$$r' = \frac{p_r}{\sqrt{n_0^2(1 - \alpha^2 r^2) - p_r^2 - \frac{p_\theta^2}{r^2}}}, \quad p'_r = \frac{1}{\sqrt{n_0^2(1 - \alpha^2 r^2) - p_r^2 - \frac{p_\theta^2}{r^2}}} \left( \frac{p_\theta^2}{r^3} - n_0^2 \alpha^2 r \right)$$

$$\theta' = \frac{1}{\sqrt{n_0^2(1 - \alpha^2 r^2) - p_r^2 - \frac{p_\theta^2}{r^2}}} \frac{p_\theta}{r}, \quad p'_\theta = 0$$

4 分, 每个方程 1 分

### 三、20%

解: (1) 触地点的坐标是

$$\vec{r}_A = (x, y, -a)^T, \quad \vec{r}_{CA} = (0, 0, -a)^T$$

$$\vec{v}_A = \vec{\omega} \times \vec{r}_{CA} + \vec{v}_C = \begin{vmatrix} \vec{e}_x & \vec{e}_y & \vec{e}_z \\ \cos \phi \dot{\theta} + \sin \phi \sin \theta \dot{\psi} & \sin \phi \dot{\theta} - \cos \phi \sin \theta \dot{\psi} & \dot{\phi} + \cos \theta \dot{\psi} \\ 0 & 0 & -a \end{vmatrix} + \begin{pmatrix} \dot{x} \\ \dot{y} \\ 0 \end{pmatrix}$$

$$= \begin{pmatrix} \dot{x} - a(\sin \phi \dot{\theta} - \cos \phi \sin \theta \dot{\psi}) \\ \dot{y} + a(\cos \phi \dot{\theta} + \sin \phi \sin \theta \dot{\psi}) \\ 0 \end{pmatrix}$$

2 分

约束方程是

$$\begin{cases} \dot{x} - a(\sin \phi \dot{\theta} - \cos \phi \sin \theta \dot{\psi}) = 0 \\ \dot{y} + a(\cos \phi \dot{\theta} + \sin \phi \sin \theta \dot{\psi}) = 0 \end{cases}$$

虚位移满足的约束方程

$$\begin{cases} \delta x - a(\sin \phi \delta \theta - \cos \phi \sin \theta \delta \psi) = 0 \\ \delta y + a(\cos \phi \delta \theta + \sin \phi \sin \theta \delta \psi) = 0 \end{cases}$$

4 分; 每个约束 2 分

(2) 拉格朗日函数是

$$L = T = \frac{1}{2} \cdot \frac{2}{5} m a^2 \cdot \vec{\omega}^2 + \frac{1}{2} m \vec{v}_C^2$$

$$= \frac{1}{5} m a^2 \left\{ (\cos \phi \dot{\theta} + \sin \phi \sin \theta \dot{\psi})^2 + (\sin \phi \dot{\theta} - \cos \phi \sin \theta \dot{\psi})^2 + (\dot{\phi} + \cos \theta \dot{\psi})^2 \right\}$$

$$+ \frac{1}{2} m (\dot{x}^2 + \dot{y}^2)$$

$$= \frac{1}{5} m a^2 \{ \dot{\phi}^2 + \dot{\theta}^2 + \dot{\psi}^2 + 2 \cos \theta \dot{\phi} \dot{\psi} \} + \frac{1}{2} m (\dot{x}^2 + \dot{y}^2)$$

6 分

平动项 2 分, 转动项 4 分

(3) 由达朗贝尔原理, 并利用约束方程消去  $\delta x, \delta y$ ,

$$m a \ddot{x} (\sin \phi \delta \theta - \cos \phi \sin \theta \delta \psi) - m a \ddot{y} (\cos \phi \delta \theta + \sin \phi \sin \theta \delta \psi)$$

$$+ \frac{2}{5} m a^2 \delta \phi \frac{d}{dt} (\dot{\phi} + \cos \theta \dot{\psi}) + \frac{2}{5} m a^2 (\ddot{\theta} + \sin \theta \dot{\phi} \dot{\psi}) \delta \theta$$

$$+ \frac{2}{5} m a^2 \delta \psi \frac{d}{dt} (\dot{\psi} + \cos \theta \dot{\phi}) = 0$$

2 分

$$\begin{cases} \frac{2}{5} m a^2 \frac{d}{dt} (\dot{\phi} + \cos \theta \dot{\psi}) = 0 \\ \frac{2}{5} m a^2 (\ddot{\theta} + \sin \theta \dot{\phi} \dot{\psi}) + m a \ddot{x} \sin \phi - m a \ddot{y} \cos \phi = 0 \\ \frac{2}{5} m a^2 \frac{d}{dt} (\dot{\psi} + \cos \theta \dot{\phi}) - m a \ddot{x} \cos \phi \sin \theta - m a \ddot{y} \sin \phi \sin \theta = 0 \end{cases}$$

6 分; 每个式子 2 分

### 四、15%

解: (1) 有多种解法。利用可积条件,

$$dp \wedge dq = d\left(\frac{Q}{P}\right) \wedge \alpha d(P^2) = 2\alpha dQ \wedge dP$$

当  $\alpha = 1/2$  是正则变换。

3 分

(2)

$$dF = pdq - PdQ = \frac{1}{2}P^2d\left(\frac{Q}{P}\right) - PdQ = -\frac{1}{2}PdQ - \frac{1}{2}QdP \Rightarrow F = -\frac{1}{2}PQ$$

2 分

$$F_2(t, q, P) = F + PQ = \frac{1}{2}PQ = \frac{1}{2}P \cdot Pq = \frac{1}{2}P^2q$$

2 分

(3)

$$\begin{aligned} K &= H + \frac{\partial F_2}{\partial t} = p(q^2 + 1) = \frac{1}{2}P^2\left(\frac{Q^2}{P^2} + 1\right) = \frac{1}{2}(P^2 + Q^2) \\ \Rightarrow \dot{Q} &= P, \quad \dot{P} = -\dot{Q} \end{aligned}$$

5 分；哈密顿量 1 分，方程各 2 分

(4)

$$Q = \pm\sqrt{2pq}, \quad P = \pm\sqrt{2p}$$

$F_1(p, Q, t)$ :  $\frac{\partial Q}{\partial p} \neq 0$ , 存在;

$F_3(p, Q, t)$ :  $\frac{\partial Q}{\partial q} \neq 0$ , 存在;

$F_4(p, P, t)$ :  $\frac{\partial P}{\partial q} = 0$ , 不存在。

3 分；每个判断 1 分

五、15%

解：(1)

$$\begin{aligned} p_r &= mr \\ p_\theta &= mr^2\dot{\theta} \\ p_\phi &= mr^2\sin^2\theta\dot{\phi} \\ H &= p_r\dot{r} + p_\theta\dot{\theta} + p_\phi\dot{\phi} - L = \frac{1}{2m}\left(p_r^2 + \frac{p_\theta^2}{r^2} + \frac{p_\phi^2}{r^2\sin^2\theta}\right) + \frac{1}{2}kr^2 \end{aligned}$$

4 分

(2)

$$\begin{aligned} \frac{d}{dt}J^2 &= [J^2, H] = \left[p_\theta^2 + p_\phi^2A(\theta), \frac{1}{2m}\left(p_r^2 + \frac{p_\theta^2}{r^2} + \frac{p_\phi^2}{r^2\sin^2\theta}\right) + \frac{1}{2}kr^2\right] \\ &= \left[p_\theta^2 + p_\phi^2A(\theta), \frac{1}{2m}\left(\frac{p_\theta^2}{r^2} + \frac{p_\phi^2}{r^2\sin^2\theta}\right)\right] \\ &= \left[p_\theta^2, \frac{1}{\sin^2\theta}\right]\frac{p_\phi^2}{2mr^2} + [A(\theta), p_\theta^2]\frac{p_\phi^2}{2mr^2} = 0 \end{aligned}$$

3 分

$$\begin{aligned} \Rightarrow \left[p_\theta^2, A(\theta) - \frac{1}{\sin^2\theta}\right] &= 0 \Rightarrow \frac{\partial}{\partial\theta}\left\{A(\theta) - \frac{1}{\sin^2\theta}\right\} = 0 \\ \Rightarrow A(\theta) &= \frac{1}{\sin^2\theta} + c, \quad c \in \mathbb{R} \end{aligned}$$

2 分

(3)

$$\begin{aligned} \delta r &= \epsilon[r, J^2] = \epsilon[r, p_\theta^2 + p_\phi^2A(\theta)] = 0 \\ \delta\theta &= \epsilon[\theta, p_\theta^2 + p_\phi^2A(\theta)] = \epsilon[\theta, p_\theta^2] = 2\epsilon p_\theta \\ \delta\phi &= \epsilon[\phi, p_\theta^2 + p_\phi^2A(\theta)] = \epsilon[\phi, p_\phi^2A(\theta)] = 2p_\phi A(\theta) \\ \delta p_r &= \epsilon[p_r, p_\theta^2 + p_\phi^2A(\theta)] = 0 \end{aligned}$$

$$\delta p_\theta = \epsilon [p_\theta, p_\theta^2 + p_\phi^2 A(\theta)] = \epsilon p_\phi^2 [p_\theta, A(\theta)] = -\epsilon p_\phi^2 \frac{\partial A(\theta)}{\partial \theta} = 2\epsilon p_\phi^2 \frac{\cos \theta}{\sin^3 \theta}$$

$$\delta p_\phi = \epsilon [p_\phi, p_\theta^2 + p_\phi^2 A(\theta)] = 0$$

6 分；每个式子 1 分

六、15%

解：设单摆的悬挂点距离平板的边缘距离为  $l$ ，则拉氏函数是

$$L = \frac{1}{2}ma^2\dot{\theta}^2 + mg\{a \sin \alpha \cos \theta + l \sin \alpha\}$$

微振动时，

$$L \approx \frac{1}{2}ma^2\dot{\theta}^2 + mg \sin \alpha \left\{ a \left( 1 - \frac{1}{2}\theta^2 \right) + l \right\}$$

$$\Leftrightarrow L = \frac{1}{2}ma^2\dot{\theta}^2 - \frac{1}{2}mga \sin \alpha \theta^2$$

6 分

利用量纲分析，绝热不变量是

$$\frac{1}{2}mga \sin \alpha \theta_M^2 \cdot 2\pi \sqrt{\frac{ma^2}{mga \sin \alpha}} \propto \theta_M^2 \sqrt{\sin \alpha}$$

6 分

$$\theta_M^2 \sqrt{\sin 90^\circ} = \theta_0^2 \sqrt{\sin 30^\circ} = \frac{\theta_0^2}{\sqrt{2}}$$

$$\theta_M = 2^{-1/4} \theta_0 \approx 0.841 \theta_0$$

3 分